# Microstructure and current-voltage characteristics of praseodymium-doped zinc oxide varistors containing MnO<sub>2</sub>, Sb<sub>2</sub>O<sub>3</sub> and Co<sub>3</sub>O<sub>4</sub>

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The effects of the oxide additives  $MnO_2$ ,  $Sb_2O_3$  and  $Co_3O_4$  commonly incorporated in commercial  $Bi_2O_3$ -doped ZnO varistors on the microstructure and the current-voltage characteristics of 0.5 mol%  $Pr_6O_{11}$ -doped ZnO varistors have been studied. A 1 mol% addition of  $Co_3O_4$  to the ZnO- $Pr_6O_{11}$  binary system does not produce any additional secondary phases in addition to Pr oxides. In contrast to this, the addition of 1 mol%  $MnO_2$  complicates the simple two-phase microstructure of the 0.5 mol%  $Pr_6O_{11}$ -doped material by forming the perovskite-type structure  $Pr(Mn_{1-x}Zn_x)O_3$ . The addition of  $Sb_2O_3$  produces the pyrochlore  $Pr_3Zn_2Sb_3O_{14}$  and spinel  $Zn_7Sb_2O_{12}$  phases. Application of the Fresnel fringe technique and the diffuse dark field technique showed that there was an amorphous Pr-rich layer of thickness <2 nm residing between adjacent ZnO grains in all the samples studied. © 2002 Kluwer Academic Publishers

### 1. Introduction

Zinc oxide ceramics containing  $Pr_6O_{11}$  and  $Co_3O_4$  and exhibiting non-ohmic current characteristics were first reported by Mukae *et al.* [1]. Pr-doped ZnO varistors have a simple two-phase microstructure of ZnO grains and a Pr oxide intergranular phase [1–3]. In comparison with Bi-doped ZnO varistors, which have other secondary phases such as spinel that play no significant electrical role, this two-phase microstructure increases the active grain boundary area through which electrical current flows, because the I-V non-linearity originates from the grain boundaries between ZnO grains.

Pr-doped ZnO varistors are not as well studied and characterised as the Bi-doped ZnO varistors. For the most part, only the effects of Co-doping on the microstructure and electrical properties have been studied [1–6], although there have been reports on grain growth control in varistors doped with both Co<sub>3</sub>O<sub>4</sub> and K<sub>2</sub>O [7–9]. The effects of other additive oxides such as MnO<sub>2</sub> and Sb<sub>2</sub>O<sub>3</sub> commonly found in the well-studied ZnO-Bi<sub>2</sub>O<sub>3</sub> system are not reported in the literature on Pr-doped ZnO varistors. MnO<sub>2</sub> is a 'performance enhancer' in Bi-doped ZnO varistors, just like Co<sub>3</sub>O<sub>4</sub>, and is added to increase the non-linear coefficient,  $\alpha$ [10], while Sb<sub>2</sub>O<sub>3</sub> is added primarily to control the ZnO grain growth [11].

In this investigation, the effects of  $MnO_2$  and  $Sb_2O_3$ on the microstructure and the electrical characteristics have been studied in a binary  $ZnO-Pr_6O_{11}$  system. X-ray powder diffraction (XRPD), scanning electron microscopy (SEM), transmission electron microscopy (TEM) and standard electrical measurement procedures were used. The purpose of this study was to ascertain whether these oxides would affect beneficially the ZnO-Pr<sub>6</sub>O<sub>11</sub> system, just as they do in the ZnO-Bi<sub>2</sub>O<sub>3</sub> system. In addition, a sample doped with Co<sub>3</sub>O<sub>4</sub> was also prepared and studied. It was used as a 'control' since the effects of Co<sub>3</sub>O<sub>4</sub> as an additive in the ZnO-Pr<sub>6</sub>O<sub>11</sub> system has been studied extensively. The possibility of a Pr-rich amorphous layer at ZnO-ZnO grain boundaries similar to those detected in Bi-doped ZnO varistors was investigated in these Pr-doped ZnO varistors using Fresnel fringe imaging and diffuse dark field imaging techniques.

# **2. Experimental procedure** 2.1. Sample preparation

High purity oxide powder starting materials were used for the preparation of the ZnO varistor samples. The powder mixtures of the selected composition were ballmilled for 24 hours using zirconia beads in deionised water. The mixtures were then dried, pressed into pellets, and sintered. The compositions chosen are summarised in Table 1. The separate effects of  $MnO_2$ ,  $Co_3O_4$  and  $Sb_2O_3$  were examined by adding them individually at 1 or 2 mol% levels to a ZnO–0.5 mol%  $Pr_6O_{11}$  binary system—samples ZPM, ZPC and ZPS. The combined effect of these additives was also

TABLE I Summary of sample compositions and material parameters obtained for  $Pr_6O_{11}\mbox{-}doped$  ZnO varistors sintered for 1 hr at  $1350^\circ C$ 

Sample	Composition	Average grain size (µm)	Non-linear coefficient, $\alpha^a$
ZP	ZnO-0.5mol%Pr <sub>6</sub> O <sub>11</sub>	16.0	3.3
ZPM	ZnO-0.5mol% Pr <sub>6</sub> O <sub>11</sub> -1mol%MnO <sub>2</sub>	15.1	9.2
ZPC	ZnO-0.5mol% Pr6O11-1mol%Co3O4	34.7	7.1
ZPS	ZnO-0.5mol% $Pr_6O_{11}-2mol\%Sb_2O_3$	3.2	5.1
ZPSMC	$\label{eq:20} \begin{split} ZnO-0.5mol\%Pr_6O_{11}-\\ 2mol\%Sb_2O_3-1mol\%\\ MnO_2-1mol\%Co_3O_4 \end{split}$	4.8	6.5

 ${}^{a}\alpha = \frac{\log(J_2/J_1)}{\log(E_2/E_1)}$ , where  $J_1 = 0.5 \text{ mA/cm}^2$  and  $J_2 = 5 \text{ mA/cm}^2$ .

demonstrated in the sample ZPSMC. All the pellets were sintered in an atmosphere of ambient air for 1 hr at  $1350^{\circ}$ C and then cooled at  $5^{\circ}$ C min<sup>-1</sup> to room temperature.

# 2.2. Characterisation

For electrical measurements, the as-sintered specimens were lapped on both surfaces to ensure flat and parallel surfaces. They were coated with conductive silver paint on both surfaces, then heat cured at 140°C for 10 min to provide ohmic contacts. The current-voltage (I-V) characteristics were determined at room temperature using a variable dc power supply (Heathkit Model IP-17). X-ray analyses of the sintered samples were carried out using Cu K<sub> $\alpha$ </sub> radiation on a Philips PW1710 vertical powder diffractometer. Microstructures were examined initially by SEM using a JSM 820 and a JEOL 5800 equipped with energy dispersive X-ray analysis (EDX). Specimens for TEM were prepared using standard ion beam thinning methods and examined using a JEOL 2000FX at 200 kV with an ultra thin Be window for EDX, attached to a Link Analyser model 6284. The microanalyses were made quantitative by using the Link Systems RTS-2/FLS computer programme. Scanning transmission electron microscopy (STEM) was also used at 100 kV using a VG-STEM to study ZnO-ZnO grain boundaries.

# 3. Results and discussion

# 3.1. X-ray powder diffraction

XRPD traces for samples containing the different additives are shown in Fig. 1. Other than the major ZnO phase, praseodymium oxides (cubic Pr<sub>6</sub>O<sub>11</sub> and hexagonal Pr<sub>2</sub>O<sub>3</sub>) are detectable in ZP, ZPC and ZPM, but not in ZPS and ZPSMC. Peaks in addition to those from the ZnO and Pr oxides were observed for samples containing MnO<sub>2</sub> (ZPM). A search of the JCPDS database did not enable the additional peaks in ZPM to be labelled. However, it has been reported separately that the phase PrMnO<sub>3</sub> occurs as a cubic perovskite phase with cell parameter a = 3.82 Å [12]. The additional peaks observed in ZPM could be indexed according to the X-ray data reported by Vickery and Klann [12]. Further analysis using EDX in SEM and TEM showed a phase containing Zn, as well as Pr and Mn. Consistent indexing of electron diffraction patterns obtained from this phase using TEM enabled the crystal lattice of this

phase to be deduced as a tetragonal *P* lattice with unit cell parameters a = 5.4 Å and c = 7.7 Å. The additional peaks observed in the XRPD patterns could be then indexed according to this crystal lattice. The lattice of this phase can be viewed as a 'cubic *C*' lattice (but without the four triads consistent with the minimum symmetry requirements for cubic symmetry) with a = 7.7 Å, which is about twice that of the PrMnO<sub>3</sub> phase. Hence, it is highly plausible that the phase we have observed is a derivative of PrMnO<sub>3</sub> with a perovskite-type structure with a chemical formula Pr(Mn<sub>1-x</sub>Zn<sub>x</sub>)O<sub>3</sub>. A detailed description of these results will be given in Section 3.2.3.

The two samples doped with Sb<sub>2</sub>O<sub>3</sub> (ZPS and ZPSMC) had a number of unidentified peaks in the XRPD traces, even after taking into account the peaks for the Zn<sub>7</sub>Sb<sub>2</sub>O<sub>12</sub> spinel phase. A phase with the correct interplanar spacings to explain these unidentified peaks was observed using TEM. From a study of its electron diffraction patterns and using EDX, this phase was identified to be a Pr-Zn-Sb pyrochlore-type phase (Zn<sub>2</sub>Sb<sub>3</sub>Pr<sub>3</sub>O<sub>14</sub>) similar to the pyrochlore Zn<sub>2</sub>Sb<sub>3</sub>Bi<sub>3</sub>O<sub>14</sub> phase formed in ZnO varistors doped with Bi<sub>2</sub>O<sub>3</sub> and Sb<sub>2</sub>O<sub>3</sub>. The crystal lattice for this phase deduced from a self-consistent indexing of the selected area electron diffraction patterns is cubic *F* with a cell parameter *a* = 10.4 Å. A detailed description of the results will be presented in Section 3.2.4.

# 3.2. Microstructural observations

Microstructural observations in the present investigation for ZP and ZPC are consistent with those reported in the literature for Pr-doped systems with or without additional Co-doping [2, 4]. Only two phases are present in the ceramics, namely, ZnO grains and an intergranular phase composed of praseodymium oxide. The microstructure of ZPM consists mainly of ZnO, Pr oxide and Pr( $Mn_{1-x}Zn_x$ )O<sub>3</sub>. The microstructure of ZPS and ZPSMC was more complex with two secondary phases, namely the spinel and the Pr-Zn-Sb pyrochloretype phases. In this respect, ZPS has a microstructure similar to ZnO varistors doped with Bi<sub>2</sub>O<sub>3</sub> and Sb<sub>2</sub>O<sub>3</sub>. Pr oxides were absent in ZPS and ZPSMC.

# 3.2.1. ZnO grains

Low magnification secondary electron images (SEIs) revealing the typical microstructure of  $Pr_6O_{11}$ -doped ZnO varistors are shown in Fig. 2. The average ZnO grain sizes shown in Table I were obtained from these micrographs by multiplying the average linear intercept length of the grains by 1.56 [13]. The grains were more uniform in samples containing Sb<sub>2</sub>O<sub>3</sub>. The ZnO grains were about 4  $\mu$ m in size for samples containing Sb<sub>2</sub>O<sub>3</sub>, and greater than 15  $\mu$ m in size for samples without Sb<sub>2</sub>O<sub>3</sub>. This grain refinement is due to the formation of the spinel phase, which, just as for Bi<sub>2</sub>O<sub>3</sub>-doped ZnO varistors [14], hinders the growth of the ZnO grains during sintering. The spinel phase forms during heating at ~900°C and is largely out of solution at the sintering temperature [14].

 $MnO_2$  and  $Co_3O_4$  exhibited different effects on ZnO grain growth. The average ZnO grain size for the



*Figure 1* X-ray powder diffraction patterns (Cu  $K_{\alpha}$ ) for Pr<sub>6</sub>O<sub>11</sub>-doped ZnO varistors sintered at 1350°C (Sp: Zn<sub>7</sub>Sb<sub>2</sub>O<sub>12</sub> spinel phase; Py: Pr-Zn-Sb pyrochlore phase).

sample containing  $Co_3O_4$  is nearly twice that containing MnO<sub>2</sub>. This is likely to be due to the formation of the secondary phase  $Pr(Mn_{1-x}Zn_x)O_3$ , during sintering, enabling it to hinder the ZnO grain growth in ZPM. In all the specimens, Mn or Co was found dissolved in the ZnO grains. From EDX analyses (Table II), the amount was about 0.2 wt% Mn and 4 wt% Co in ZPM and ZPC respectively. Thus, while some Mn does dissolve in ZnO grains, virtually all of the Mn added precipitates out in the  $Pr(Mn_{1-x}Zn_x)O_3$  secondary phase, whereas all the Co dissolves into the ZnO grains. The dissolution of Mn and Co in the ZnO grains, even if relatively small as in the case of Mn, is important for the electrical properties of varistors, because these two elements act as dopants, as discussed in Section 3.3.

In all the cases examined, there was no dissolution of Pr in the ZnO grains. This is in agreement with other work [6, 15]. The  $Pr^{3+}$  cations are considerably larger than  $Zn^{2+}$  cations and solid solution between Pr oxide and ZnO is severely limited. This explains why Pr mainly segregates to the grain boundary regions, producing either thin layers of Pr-rich material at the ZnO grain boundaries or second phase Pr oxides, and hence giving rise to the varistor effect.

#### 3.2.2. Praseodymium oxides

TEM observations showed that  $Pr_6O_{11}$  particles existed in ZP, ZPM and ZPC, as expected from the XRPD results. The identity of the particles was confirmed using EDX and also by indexing selected area electron diffraction patterns (SADPs) obtained from the phase. The particles existed either as a single crystalline grain or a polycrystalline aggregate, with grain sizes ranging from 0.5 to 2  $\mu$ m.  $Pr_6O_{11}$  particles were

TABLE II Average compositions (wt% of metallic elements) together with standard deviations obtained using TEM/EDX analyses of ZnO grains in ZPM, ZPC and ZPSMC; spinel grains and Pr-Zn-Sb pyrochloretype grains in ZPS and ZPSMC. At least 5 grains of each type were sampled for each specimen

Elements (wt%)	ZPM	ZPC	ZPS	ZPSMC
ZnO grains				
Zn	$99.8\pm0.1$	$96.2\pm0.7$		$96.4\pm0.2$
Mn	$0.2\pm0.1$	-		$0.6\pm0.1$
Co	-	$3.8\pm0.7$		$2.7\pm0.2$
Spinel grains				
Zn			$63.6\pm0.5$	$58.9 \pm 1.9$
Sb			$34.8\pm1.5$	$33.7\pm1.5$
Mn			-	$2.0\pm0.4$
Co			-	$5.4\pm0.8$
Pr-Zn-Sb pyrochlore-				
type phase				
Zn			$18.6\pm0.4$	$18.5\pm3.0$
Sb			$38.1\pm0.5$	$35.7 \pm 1.5$
Pr			$43.3\pm0.6$	$43.0\pm2.2$
Mn			_	$1.0\pm0.2$
Co			-	$1.8\pm0.1$

(a) ZP



(c) ZPC



(e) ZPSMC



Figure 2 SEM micrographs of the Pr<sub>6</sub>O<sub>11</sub>-doped ZnO samples.

found at grain boundaries, triple point junctions or embedded in ZnO grains.

Although  $Pr_2O_3$  was detected in ZP, ZPM and ZPC using X-ray diffraction, TEM observations did not show the existence of this Pr oxide phase. However, the amount of  $Pr_2O_3$  is significantly lower than  $Pr_6O_{11}$  in all three samples, and is particularly low in ZPM and ZPC, so one explanation for this failure to observe  $Pr_2O_3$  particles is the relatively small number of Pr oxide grains (<20) sampled during the TEM work.

It was apparent from the backscattered electron images (BEIs) for ZPC, such as the one shown in Fig. 3, that in parts of the sample there are regions where there is a virtually continuous thick Pr oxide intergranular layer between ZnO grains which can be readily identified in the SEM. Other workers have also observed such a continuous intergranular layer in Pr-doped ZnO varistors containing Co [6]. Significantly, only ZPC showed this—similar regions of continuous oxide layers were not observed in the other samples investigated here.

(b) ZPM







*Figure 3* SEM micrograph (BEI) of sample ZPC, indicating an almost continuous Pr-rich intergranular layer.

## 3.2.3. The ZnO-Pr<sub>6</sub>O<sub>11</sub>-MnO<sub>2</sub> system

A general microstructure of sample ZPM is shown in Fig. 4a. Small grains, about 0.5  $\mu$ m in size, were usually found either at triple point junctions or embedded in ZnO grains. EDX spectra, such as the example in Fig. 4b, showed the presence of Pr as well as Zn and Mn. It should be noted that the Zn peaks also contained contributions from the surrounding ZnO grains because of beam spreading. Electron diffraction patterns were also obtained, and a systematic approach was used to deduce the crystal structures for this Pr-Zn-Mn phase so that consistent indexing of these diffraction patterns was achieved. As a first attempt, the observation of centred lattices in reciprocal space suggested that the crystal lattice of this phase is base-centred, and it was found that a cubic C 'lattice' with a = 7.7 Å could index all the diffraction patterns obtained from this phase. This is in reality a tetragonal P lattice with a = b = 5.4 Å  $(=c/\sqrt{2})$  and c=7.7 Å. Electron diffraction patterns of this Pr-Zn-Mn phase are shown in Fig. 5, together with the corresponding tilt angles required to go from one zone to another.

It was suggested in Section 3.1 that the Pr-Zn-Mn phase could be a derivative of the cubic perovskite  $PrMnO_3$  phase. Considered as a cubic C 'lattice', the unit cell parameter is nearly twice that of the cubic perovskite. Most perovskite-type oxides (ABO<sub>3</sub>) involve slight modifications that are strongly related to the basic cubic symmetry. The essential requirement for the stability of the perovskite structure is that the A and B ions should be of suitable size to form 12-fold and 6fold coordination polyhedra respectively, with the ionic radii  $r_A > 0.9$  Å and  $r_B > 0.51$  Å [16]. It is known that distortion in the perovskite-type structure can lead to a larger orthorhombic unit cell. A large number of these orthorhombic perovskites, especially those featuring trivalent A and B ions have a structure typified by that of GdFeO<sub>3</sub>. The cell parameters a and b are about  $\sqrt{2}a_o$ and  $c \approx 2a_o$  where  $a_o$  is the side of the pseudo-cubic perovskite cell [16]. Complex perovskite-type oxides (A'A"BO<sub>3</sub>) exhibiting the GdFeO<sub>3</sub> structure have been reported in the literature, for example, La<sub>0.7</sub>Ca<sub>0.3</sub>MnO<sub>3</sub> [17] and  $Pr_{0.75}Sr_{0.05}MnO_3$  [18].



*Figure 4*  $Pr(Mn_{1-x}Zn_x)O_3$  phase observed in the ZnO-Pr<sub>6</sub>O<sub>11</sub>-MnO<sub>2</sub> system: (a) general microstructure; (b) Typical EDX spectrum of the  $Pr(Mn_{1-x}Zn_x)O_3$  phase.

The lattice parameters for the Pr-Zn-Mn phase deduced from electron diffraction patterns are similar to that for GdFeO<sub>3</sub>. It should be noted that the crystal structure deduced from electron diffraction patterns will tend to be either the same symmetry or, more likely, a higher symmetry than the real crystal which might be pseudo-cubic, that is, it can be slightly orthorhombic, rhombohedral or monoclinic. Such small changes in cell parameters are not readily detectable in electron diffraction patterns. The ionic radii of Zn<sup>2+</sup>, Mn<sup>3+</sup> and Mn<sup>4+</sup> for 6-fold coordination are 0.74 Å, 0.66 Å and 0.60 Å respectively [19]. The distortion of the cubic perovskite PrMnO<sub>3</sub> structure is due to the replacement of the  $Mn^{3+}$  ions by the  $Zn^{2+}$  ions, with some of the manganese ions existing as  $Mn^{4+}$  to maintain charge neutrality. This mixed valence Mn(III)-Mn(IV) of manganese has been reported in praseodyme manganites,  $Pr_{0.7}Ca_{0.25}Sr_{0.05}MnO_3$  and  $Pr_{0.75}Sr_{0.05}MnO_3$  [18].



Figure 5 Electron diffraction patterns of the  $Pr(Mn_{1-x}Zn_x)O_3$  phase tilted from one zone to another.  $\theta$  is the relevant tilt angle.

EDX analysis showed that the apparent amount of Zn detected varied over a wide range for the different Pr-Mn-Zn grains analysed. This is most probably due to the different levels of contributions from surrounding ZnO grains arising from beam spreading. However, if the Zn peak was ignored, it was found that the atomic ratio Pr : Mn was consistently about 4 : 3. This implies that 1/4 of the Mn ions are replaced by Zn<sup>2+</sup> ions. Since the crystal structure is 'cubic C' as deduced from the electron diffraction patterns, a possible crystal structure

of the  $Pr(Mn_{1-x}Zn_x)O_3$  phase can be inferred consistent with this. The structure is shown in Fig. 6, with some of the  $Mn^{3+}$  ions being replaced by  $Zn^{2+}$  ions to obtain a base-centred lattice. The positions of the oxygen ions are only shown in the smaller cubic perovskite unit cell indicated by the solid lines.

Instead of the 'ordered' structure suggested in Fig. 6, it is also relevant to consider a disordered structure due to octahedral tilting. The theory of the doubling of certain cell axes through octahedral tilting in perovskites



*Figure 6* Proposed crystal structure of  $Pr(Mn_{1-x}Zn_x)O_3$  based on the perovskite structure. Bold lines indicate unit cell, while the solid lines indicate the cubic perovskite unit cell.

was developed by Glazer [20, 21]. Recently, Thomas [22, 23] proposed a new method for the quantitative description of all perovskite ABO<sub>3</sub> structures involving the relationship between polyhedral volume ratio  $V_A/V_B$  and octahedral tilt angles. As the polyhedral volume ratio  $V_A/V_B$  decreases, the degree of octahedral tilting increases, causing a doubling of the unit cell. By replacing the smaller Mn<sup>3+</sup> ion with the bigger Zn<sup>2+</sup> ion,  $V_B$  decreases, causing the doubling of the unit cell. The full structural description of this Pr(Mn<sub>1-x</sub>Zn<sub>x</sub>)O<sub>3</sub> phase will require further work involving single crystals and characterising them using techniques such as XRD and neutron diffraction, which is beyond the scope of this report.

Since  $Pr(Mn_{1-x}Zn_x)O_3$  is clearly a derivative of the perovskite  $PrMnO_3$  phase, its formation should be similar to that of  $PrMnO_3$ . Vickery and Klann [12] followed the development of  $PrMnO_3$  through the firing cycle by XRD. X-ray evidence showed very little reaction at 900°C, but after a 2 hr treatment at 1200°C, the perovskite pattern could be identified. Hence, it is reasonable to deduce that the  $Pr(Mn_{1-x}Zn_x)O_3$  phase is formed at around 1200°C as well. Furthermore, from TEM, the morphology of the  $Pr(Mn_{1-x}Zn_x)O_3$  phase consisted of well-developed and rounded grains (e.g., Fig. 4), suggesting that these grains have had plenty of time to adjust their shape freely to a near-equilibrium form, and implying that they are most probably formed during sintering.

### 3.2.4. Samples containing Sb<sub>2</sub>O<sub>3</sub>

The general microstructure of samples containing  $Sb_2O_3$  is shown in Fig. 7a. Small grains ranging in size from 0.5 to 1.5  $\mu$ m were often observed embedded in ZnO grains or existing at triple junctions. EDX spectra of this phase (e.g., Fig. 7b) showed the presence of Pr as well as Zn and Sb. Electron diffraction patterns were also obtained from these grains, and a systematic approach was used to deduce the crystal structure for this Pr-Zn-Sb phase to achieve consistent indexing of these diffraction patterns. It was found that a cubic *F* lattice with a = 10.4 Å can index all the diffraction patterns obtained from this phase. Examples of diffraction patterns of this Pr-Zn-Sb phase are shown in Fig. 8, together with the corresponding tilt angles required to go from one zone to another.

A cubic *F* lattice with a = 10.4 Å is characteristic of a pyrochlore structure. A Bi-Zn-Sb pyrochlore phase is often found in ZnO varistors doped with Bi<sub>2</sub>O<sub>3</sub> and Sb<sub>2</sub>O<sub>3</sub>. Wong [24] defined this pyrochlore phase as an oxygen-deficient pyrochlore with formula Bi<sub>2</sub>Zn<sub>4/3</sub>Sb<sub>2/3</sub>O<sub>6</sub>. Inada [25, 26] suggested the formula Bi<sub>3</sub>Zn<sub>2</sub>Sb<sub>3</sub>O<sub>14</sub> with the lattice constant a = 10.45 Å. Hence, it is highly possible that the Pr-Zn-Sb phase observed in the ZPS and ZPSMC systems is a pyrochlore-type phase.

The composition determined from EDX microanalyses was 43.3 wt% Pr, 18.6 wt% Zn and 38.1 wt% Sb. This corresponds to 34.6 at.% Pr, 32.1 at.% Zn and 33.3 at.% Sb, which gives a Pr : Sb atomic ratio of approximately 1 : 1. Allowing for a contribution to the Zn signal from the surrounding ZnO grains, this atomic





*Figure 7*  $Pr_3Zn_2Sb_3O_{14}$  pyrochlore-type phase observed in samples containing  $Sb_2O_3$ : (a) general microstructure; (b) Typical EDX spectrum of the  $Pr_3Zn_2Sb_3O_{14}$  pyrochlore-type phase.

ratio suggests a chemical formula  $Pr_3Zn_2Sb_3O_{14}$  which is analogous to that of the  $Bi_3Zn_2Sb_3O_{14}$  pyrochlore phase suggested by Inada [25, 26]. The ionic radius of  $Pr^{3+}$  is 1.06 Å, and so it can readily substitute for the  $Bi^{3+}$  ion of radius 0.96 Å in the pyrochlore structure.

The morphology of  $Pr_3Zn_2Sb_3O_{14}$  consisted of welldeveloped grains. In  $Bi_2O_3$ -doped ZnO varistors containing  $Sb_2O_3$ ,  $Bi_3Zn_2Sb_3O_{14}$  pyrochlore is formed during cooling and arises from a chemical reaction between liquid  $Bi_2O_3$  and the  $Zn_7Sb_3O_{12}$  spinel phase [25, 26]. Liquid  $Pr_6O_{11}$  does not exist during the sintering of Pr-doped ZnO varistors because of its high melting point, and so the precise formation mechanism of  $Pr_3Zn_2Sb_3O_{14}$  will be different from that of  $Bi_3Zn_2Sb_3O_{14}$ . A detailed analysis of this mechanism needs to be performed, but is outside the scope of the present work.



*Figure 8* Electron diffraction patterns of  $Pr_3Zn_2Sb_3O_{14}$  pyrochlore-type phase tilted from one zone to another.  $\theta$  is the tilt angle required to go from one zone to another.

## 3.2.5. Intergranular microstructure at ZnO-ZnO grain boundaries

For samples not containing  $Sb_2O_3$ , i.e., ZP, ZPM and ZPC, a high proportion of the grain boundaries between adjacent ZnO grains contained clearly observable thin (<2 nm thick) continuous second phase films with dark contrast attributable to the presence of heavy scattering elements, and therefore consistent with the presence of Pr. Invariably, such intergranular material is connected to a triple junction.

Examination of these ZnO grain boundary films by diffuse dark field imaging indicated that they were amorphous. Through-focal series of Fresnel images from such grain boundaries aligned parallel to the electron beam were also performed which further confirmed the presence of a thin boundary layer. A typical example is shown in Fig. 9 for ZPM. The boundaries exhibit a bright central fringe underfocus and a dark central fringe overfocus, which is characteristic of a lower scattering potential at the boundary relative to the surrounding bulk material. The thickness of the boundary ranged from 0.4 to 1.9 nm. Since the Fresnel technique is known to overestimate the thickness of the layer by 20% to 35% compared to those determined by HREM, which are closer to the actual widths [27], the actual thickness of these boundaries should range from 0.3 to 1.6 nm.

For samples containing  $Sb_2O_3$ , the presence of an intergranular phase was also obvious from the darker contrast at the grain boundaries in conventional bright field imaging. However, intriguingly, Fresnel fringes could not be observed when a through-focal series were performed, presumably because the contrast between

the bright and dark fringes was too small to be observable. Instead, further analysis was conducted in the VG-STEM at 100 kV, to obtain the composition of this intergranular layer in sample ZPSMC. Line scans for each element across the boundaries were obtained, as shown in Fig. 10. These line scans indicated that very little or no Zn is dissolved in the Pr-rich intergranular phase, but clear segregation of Pr and Mn and a suggestion of some segregation of Co was apparent. It is likely that these species also segregated at the ZnO grain boundaries in the other ZnO-Pr<sub>6</sub>O<sub>11</sub> systems containing the individual additives, i.e., in samples ZPM, ZPC and ZPS.

There are only limited reports in the literature on the nature of the grain boundary regions of Pr-doped ZnO varistors [6, 28]. Solórzano et al. [6] found that the majority of the grain boundaries in a Co-doped and Pr-doped ZnO varistor appeared precipitate-free, and that these grain boundaries exhibited Pr segregation. In contrast to this, Co was found to be distributed homogeneously in the ZnO matrix. Lee et al. [28] also observed a similar result in their Co-doped and Prdoped ZnO varistors. Solórzano et al. [6] also used HREM to observe a precipitate-free ZnO-ZnO grain boundary. They observed that the lattice fringes of the adjacent ZnO grains were continuous right up to the plane of this particular boundary, consistent with the absence of any amorphous layer at the boundary plane. Early reports on the ZnO-Bi<sub>2</sub>O<sub>3</sub> varistor system were inconsistent regarding the observation of a similar Bi-rich amorphous phase. However, recent reports [29, 30] show that such a thin amorphous layer is



*Figure 9* A typical boundary in sample ZPC observed by through-focal series of Fresnel images, (a) underfocus, (b) overfocus, (c) diffuse dark field image, and (d) plot of the distance between fringe maxima as a function of defocus.

a common feature. Our own work here suggests that this too occurs in Pr-doped ZnO varistors, and that there is an amorphous Pr-rich layer <2 nm thick residing between ZnO grains. However, further work is required, preferably using HREM as well as Fresnel fringe analysis, and examining a number of boundaries, in order to give results that are statistically robust.

# 3.3. I-V characteristics

The electrical properties of the  $Pr_6O_{11}$ -based ZnO materials were characterised by their electric field-current density (E-J) properties. Typical E-J curves are shown in Fig. 11. The corresponding non-linear coefficients are summarised in Table I. The non-linear coefficients,  $\alpha$ , were determined between the current density range of 0.5 mA/cm<sup>2</sup> and 5 mA/cm<sup>2</sup>.



Figure 10 Digital line scans of the different elements in ZPSMC. The grain boundary is situated at the '53' pixel position. 1 pixel = 0.5 nm.

It is evident from the curves that the resistivity of the varistors is increased markedly in samples ZPC, ZPS and ZPSMC. These three samples also have significantly higher threshold electrical field strengths before non-linearity begins. Of these three samples, the change from linear I-V characteristics to non-linear characteristics is most abrupt for ZPC, which also has the highest  $\alpha$  of these three samples. Interestingly, ZPM has a higher  $\alpha$  than ZPC, but this is at the expense of being a superior conductor prior to non-linearity, having a gentle transition from linearity to non-linearity, and a threshold electrical field strength over an order of magnitude lower than ZPC. It is therefore apparent that Co is, by

some margin, the first additional dopant of choice for Pr-doped ZnO varistors out of Co, Mn and Sb.

Mukae *et al.* [1] suggested that the origin of the varistor action in ZnO ceramics doped with  $Pr_6O_{11}$  and  $Co_3O_4$  can be attributed to the formation of electrical depletion layers at the interfacial regions between the almost continuous Pr oxide intergranular layer and the Co-doped ZnO grains. Electrons trapped by deep surface states at the grain boundaries form two Schottky barriers on both sides of the ZnO grain boundaries, just as in the case of Bi<sub>2</sub>O<sub>3</sub>-doped ZnO varistors. However, in our work, the formation of a partially continuous thick Pr oxide intergranular phase was only observed



Figure 11 Electric field-current density (E-J) curves for the  $Pr_6O_{11}$ -doped ZnO varistors.

in the system containing  $Co_3O_4$  (ZPC). It was not found in the other systems containing  $MnO_2$  and/or  $Sb_2O_3$ : ZPM, ZPS and ZPSMC. This clearly suggests that it is not necessary to have a thick continuous intergranular Pr oxide layer to produce varistor behaviour. Indeed, studies on Bi<sub>2</sub>O<sub>3</sub>-doped ZnO varistors have shown that simply coating the ZnO grains with a Bi-rich phase is not sufficient to produce good varistor electrical properties, and may even be detrimental [33].

At present, as we have stated elsewhere [34], there is no unique model that can predict quantitatively ZnO varistor I-V behaviour from a knowledge of the microstructure. Only a qualitative explanation is possible. Levinson and Philipp [35] have pointed out that, in practice, there are a variety of inter-grain conduction paths that operate in parallel in varistors. These can be through the grain boundary region or through the bulk intergranular material, and are sensitive to the presence of chemical additives. It has been postulated that transition metal oxides, such as Co and Mn, are involved in the formation of interfacial states and deep bulk traps at grain boundaries [11], providing large potential barriers to give better non-linear characteristics [36]. In the varistors we have examined in this investigation, conduction can occur through the Pr-rich intergranular layers containing different chemical additives, as well as through the minor secondary phases such as perovskite  $Pr(Mn_{1-x}Zn_x)O_3$ , pyrochlore  $Pr_3Zn_2Sb_3O_{14}$ and Zn<sub>7</sub>Sb<sub>2</sub>O<sub>12</sub> spinel residing at grain boundaries and triple junctions.

The addition of  $Sb_2O_3$  in the ZnO-Bi<sub>2</sub>O<sub>3</sub> system has the effect of enhancing the solubility of Zn ions in the Bi<sub>2</sub>O<sub>3</sub>-rich liquid phase, and so it is important for the defect distribution formed at the grain boundaries during cooling [11]. In this investigation on the ZnO-Pr<sub>6</sub>O<sub>11</sub> system, the non-linear coefficient was increased from 3 to 5 after the addition of  $Sb_2O_3$  (Table I). However, a Pr-rich liquid phase is not formed at 1350°C, so that the effect of  $Sb_2O_3$  on the non-linear behaviour is different from that which arises in Bi-doped ZnO. A recent review [37] pointed out that the Zn<sub>7</sub>Sb<sub>2</sub>O<sub>12</sub> spinel phase itself could exhibit varistor behaviour under appropriate oxidation conditions, providing an alternative electrical pathway to the grain boundary regions with an appropriate potential barrier. Nevertheless, our results from ZPS suggest that the improvement in the varistor behaviour is less significant than the addition of Co alone. The grain refinement in ZPS and ZPSMC in comparison with ZPC means that the threshold voltage per boundary for non-linear current-voltage behaviour to commence is much lower for ZPS and ZPSMC (of the order of 0.3 V and 0.6 V per boundary, simply dividing the electrical field strength by the number of grains per unit length in the sample) than for ZPC (2.4 V per boundary). By comparison, the threshold voltage for grain boundaries in ZPM is estimated to be 0.03 V per boundary.

Finally, we note that the occurrence of thin Bi-rich intergranular films in Bi-doped ZnO material has been interpreted as enhancing the varistor effect through providing a path for rapid diffusion and the formation of electrically active grain boundary states [29, 38]. Our observation of Pr-rich intergranular amorphous layers may also indicate that a similar effect arises in Pr-doped ZnO varistors.

#### 4. Conclusions

The addition of MnO<sub>2</sub> to the ZnO-Pr<sub>6</sub>O<sub>11</sub> binary system has a markedly different effect compared to the  $ZnO-Bi_2O_3$  binary system. In the  $ZnO-Bi_2O_3$  system, no secondary phases associated with MnO<sub>2</sub> are formed. However, the addition of MnO<sub>2</sub> to the ZnO-Pr<sub>6</sub>O<sub>11</sub> binary system complicates the simple two-phase microstructure by forming the perovskite-type structure  $Pr(Mn_{1-x}Zn_x)O_3$ . It improves the variator behaviour of Pr-doped ZnO material, but not as significantly as Co, the first additional dopant of choice in commercial Pr-doped ZnO varistor materials. The addition of Sb<sub>2</sub>O<sub>3</sub> produces two additional secondary phases, namely the pyrochlore Pr<sub>3</sub>Zn<sub>2</sub>Sb<sub>3</sub>O<sub>14</sub> and spinel Zn<sub>7</sub>Sb<sub>2</sub>O<sub>12</sub> phases. ZnO grain growth is inhibited with the addition of Sb<sub>2</sub>O<sub>3</sub> because of the formation of the spinel phase during sintering. Analyses of ZnO grain boundaries using Fresnel fringe imaging and digital line scans have shown that there is an amorphous Pr-rich layer <2 nm thick at these grain boundaries.

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#### References

- 1. K. MUKAE, K. TSUDA and I. NAGASAWA, Jpn. J. Appl. Phys. 16 (1977) 1361.
- 2. K. MUKAE, Ceram. Bull. 66 (1987) 1329.
- 3. K. MUKAE, K. TSUDA and S. SHIZA, *IEEE Trans Power* Delivery 3 (1988) 591.
- 4. A. B. ALLES and V. L. BURDICK, J. Appl. Phys. 70 (1991) 6883.
- 5. A. B. ALLES, R. PUSKAS, G. CALLAHAN and V. L. BURDICK, *J. Amer. Ceram. Soc.* **76** (1993) 2098.
- I. G. SOLÓRZANO, J. B. VANDER SANDE, K. K. BAEK and H. L. TULLER, in Proceedings of the 51st Annual Meeting of the Microscopy Society of America, 1993, edited by G. W. Bailey and C. L. Rieder, p. 1150.
- 7. H. K. VARMA, K. P. KUMAR, K. G. K. WARRIER and A. D. DAMODARAN, J. Mater. Sci. Lett. 8 (1989) 974.

- 8. P. M. VIPIN, V. V. SANJAYNATH, H. K. VARMA, K. G. K. WARRIER and A. D. DAMODARAN, *Br. Ceram. Trans. J.* **89** (1990) 87.
- 9. Idem., Ceram. Inter. 16 (1990) 319.
- 10. T. K. GUPTA, J. Amer. Ceram. Soc. 73 (1990) 1817.
- 11. K. EDA, IEEE Electrical Insulation Magazine 5 (1989) 28.
- 12. R. C. VICKERY and A. KLANN, J. Chem. Phys. 27 (1957) 1161.
- 13. M. I. MENDELSON, J. Amer. Ceram. Soc. 52 (1969) 443.
- 14. E. OLSSON, L. K. L. FALK, G. L. DUNLOP and R. ÖSTERLUND, J. Mater. Sci. 20 (1985) 4091.
- 15. S. Y. CHUN, N. WAKIYA, H. FUNAKUBO, K. SHINOZAKI and N. MIZUTANI, J. Amer. Ceram. Soc. 80 (1997) 995.
- 16. N. RAMADASS, Mater. Sci. Eng. 36 (1978) 231.
- Y. H. LI,K. A. THOMAS, P. S. I. P. N. DE SILVA, L. F. COHEN, A. GOYAL, M. RAJESWARI, N. D. MATHUR, M. G. BLAMIRE, J. E. EVETTS, T. VENKATESAN and J. L. MACMANUS-DRISCOLL, J. Mater. Res. 13 (1998) 2161.
- 18. M. HERVIEU, G. VAN TENDELOO, C. A. MAIGNAN and B. RAVEAU, *Phys. Rev. B* **53** (1996) 14274.
- M. H. BATTEY, in "Mineralogy for Students," 2nd ed. (Longman, New York, 1986) p. 13.
- 20. A. M. GLAZER, Acta Crystallographica B28 (1972) 3384.
- 21. Idem., ibid. A31 (1975) 756.
- 22. N. W. THOMAS, ibid. B52 (1996) 16.
- 23. Idem., ibid. B54 (1998) 585.
- 24. J. WONG, J. Appl. Phys. 46 (1975) 1653.
- 25. M. INADA, Jpn. J. Appl. Phys. 17 (1978) 1.

- 26. Idem., ibid. 19 (1980) 409.
- 27. M. K. CINIBULK, H. KLEEBE and M. RÜHLE, J. Amer. Ceram. Soc. **76** (1993) 426.
- 28. Y. S. LEE, K. S. LIAO and T. Y. TSENG, *ibid.* **79** (1996) 2379.
- 29. Y. M. CHIANG, H. WANG and J. R. LEE, J. Microscopy 191 (1998) 275.
- 30. K. KOBAYASHI, O. WADA, M. KOBAYASHI and Y. TAKADA, *J. Amer. Ceram. Soc.* **81** (1998) 2071.
- 31. A. MIRALLES, A. CORNET, A. HERMS and J. R. MORANTE, *Mater. Sci. Eng.* A109 (1989) 201.
- 32. Y. C. CHEN, C. Y. SHEN, H. Z. CHEN, Y. F. WEI and L. WU, Jpn. J. Appl. Phys. 30 (1991) 84.
- 33. J. P. GAMBINO, W. D. KINGERY, G. E. PIKE, L. M. LEVINSON and H. R. PHILIPP, J. Amer. Ceram. Soc. 72 (1989) 642.
- 34. H. H. HNG and K. M. KNOWLES, *ibid.* **83** (2000) 2455.
- 35. L. M. LEVINSON and H. R. PHILIPP, *Am. Ceram. Soc. Bull.* **65** (1986) 639.
- 36. S. EZHILVALAVAN and T. R. N. KUTTY, J. Mater. Sci. : Mater. in Elec. 7 (1996) 137.
- 37. D. R. CLARKE, J. Amer. Ceram. Soc. 82 (1999) 485.
- 38. J. R. LEE, Y. M. CHIANG and G. CEDER, Acta Mater. 45 (1997) 1247.

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